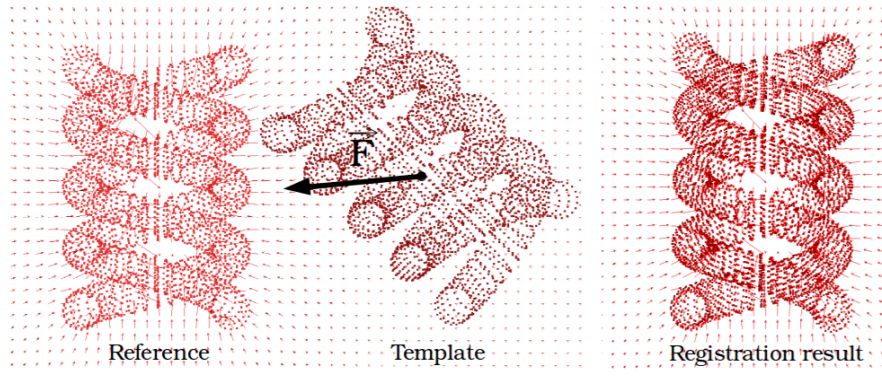


## GRAVITATIONAL METHODS FOR POINT SET ALIGNMENT

$$E(\mathbf{R}, \mathbf{t}) = -G \sum_{i,j} \frac{m_i m_j}{\|\mathbf{R}\mathbf{y}_i + \mathbf{t} - \mathbf{x}_j\|_2 + \epsilon}$$

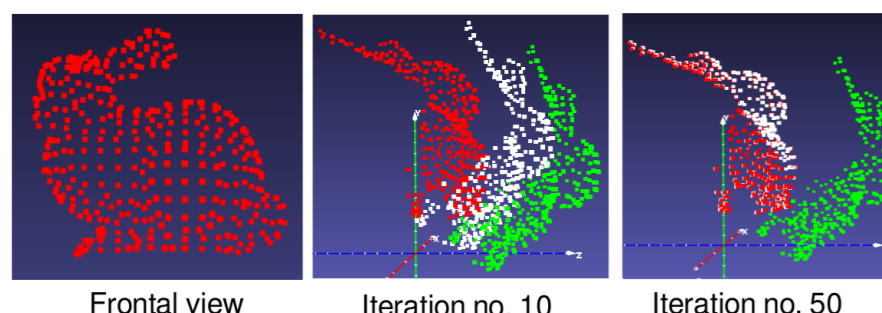
$$\tilde{\mathbf{f}}_i = -G m_j \sum_j m_j (\|\mathbf{y}_i - \mathbf{x}_j\|_2 + \epsilon)^{-3/2} \hat{\mathbf{n}}_{i,j} - \eta \mathbf{v}_i$$



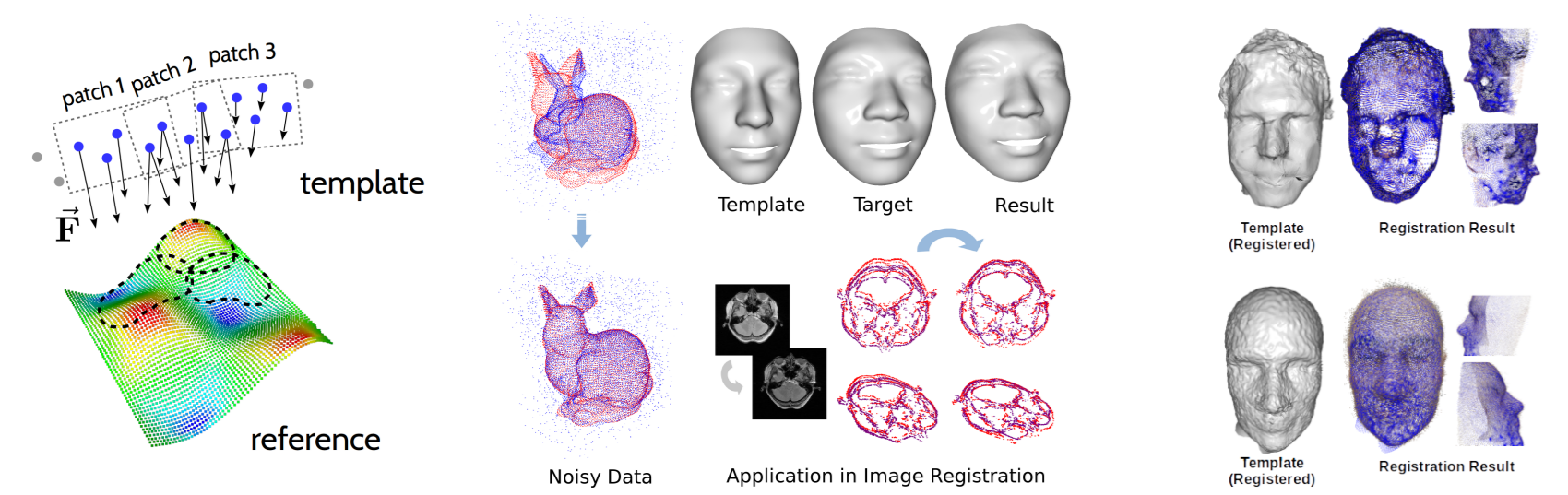
Rigid Gravitational Approach (Second-Order ODEs) [3]

$$\mathbf{F}^{Y^i} \propto g(s^{X_j, Y^i})$$

$$\mathbf{F}^{Y^i} = -G \sum_{j=1}^N g(s^{X_j, Y^i}) h(\|r^{Y^i} - r^{X_j}\|) \mathbf{n}_{i,j}$$



Modified Rigid Gravitational Approach with Shape Descriptors [4]

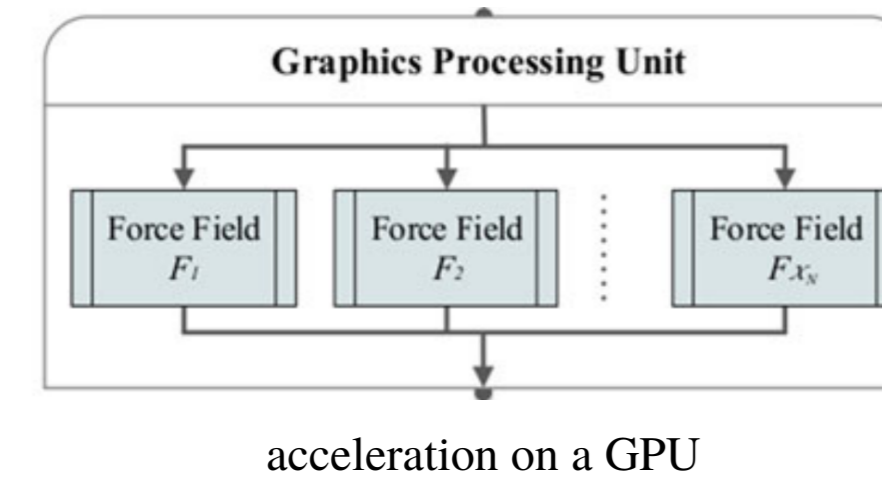


Non-Rigid Gravitational Approach (Second-Order ODEs) [5]

$$F_i^{(ex)} = -\gamma \sum_{j=1}^M \frac{m_i m_j}{\|q_i - p_j\|^3} (q_i - p_j)$$

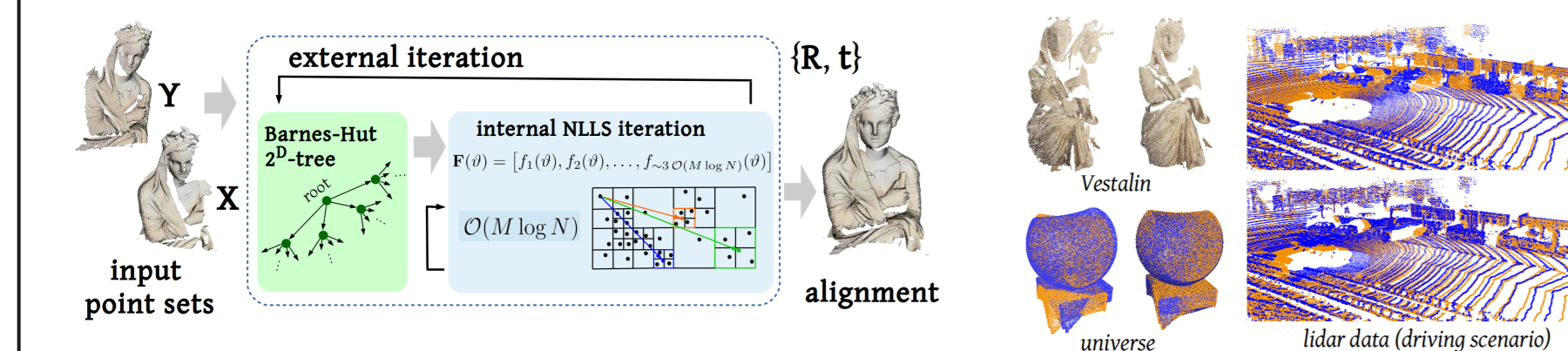
$$F_i^{(in)} = \kappa \sum_{j=1}^M \Delta_i \frac{(q_i - p_j)}{\|q_i - p_j\|^3}$$

$$\Delta_i = 1 - \frac{\|f_i^t - f_j^t\|}{\sqrt{D}} \quad \Delta_i = 2 \left( 0.5 - \frac{\|f_i^t - f_j^t\|}{\sqrt{D}} \right)$$



acceleration on a GPU

Feature-Enhanced Physics-Based Approach with Rigid Body Dynamics (Physical Heuristics and Monte-Carlo Simulation) [6]



$$\zeta^-(E(\mathbf{R}, \mathbf{t})) = \sum_{i,j} \frac{1}{G m_i m_j} \|\mathbf{R}\mathbf{y}_i + \mathbf{t} - \mathbf{x}_j\|_2 \quad \zeta^+(E(\mathbf{R}, \mathbf{t})) = \sum_{i,j} m_i m_j \|\mathbf{y}_i - \mathbf{x}_j\|_2 \approx m_{y_i} \left[ \sum_{k=1}^K m_{x_k} \right] \|\mathbf{y}_i - \bar{\mathbf{x}}\|_2$$

$$E(\mathbf{R}, \mathbf{t}) = \sum_i \sum_j m_i m_j \|\mathbf{R}\mathbf{y}_i + \mathbf{t} - \mathbf{x}_j\|_2 \quad E(\mathbf{R}, \mathbf{t}) = \sum_{y_i} m_{y_i} \sum_{k_j \in \mathcal{K}(y_i)} m_{k_j} \|\mathbf{R}\mathbf{y}_i + \mathbf{t} - \mathbf{k}_j\|_2$$

Rigid Gravitational Approach (Altered Physical Laws) [8]

## ANALYSIS OF THE ENERGY LANDSCAPE WITH CALCULUS

### Unit Circle $S^1$

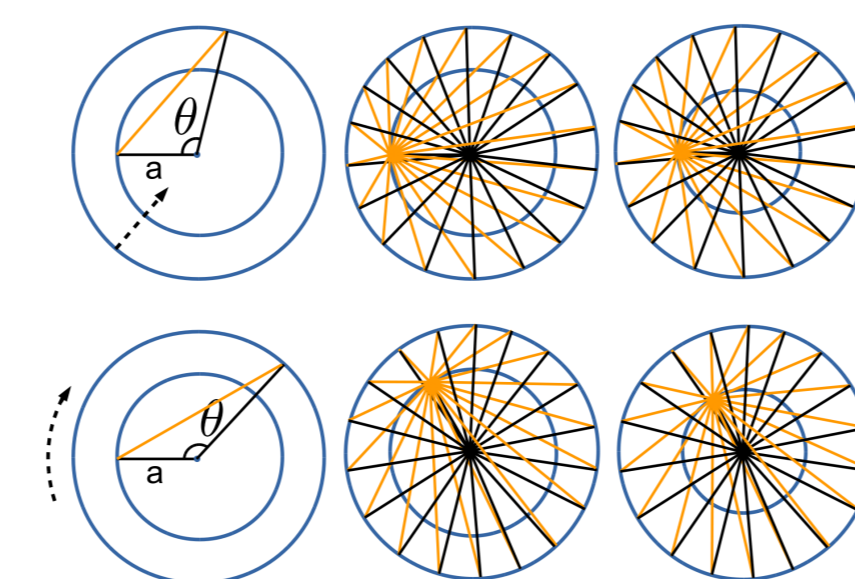
$$E_s = rNM = NM$$

$$d = \sqrt{2(1 - \cos \theta)} = 2 \sin \frac{\theta}{2}$$

$$\bar{d} = \frac{1}{\pi} \int_0^\pi 2 \sin \frac{\theta}{2} d\theta = -\frac{4}{\pi} \left( \cos \frac{\theta}{2} \right) \Big|_0^\pi = \frac{4}{\pi}$$

$$E_a = \bar{d}NM = \frac{4NM}{\pi}$$

$$\bar{d}(a) = \frac{1}{\pi} \int_0^\pi \sqrt{a^2 + 1 - 2a \cos \theta} d\theta = \begin{cases} \frac{4}{\pi} & \text{if } a = 1, \\ \frac{2|a-1|}{\pi} E\left(\frac{\theta}{2} \mid -\frac{4a}{(a-1)^2}\right) & \text{if } a \neq 1. \end{cases}$$



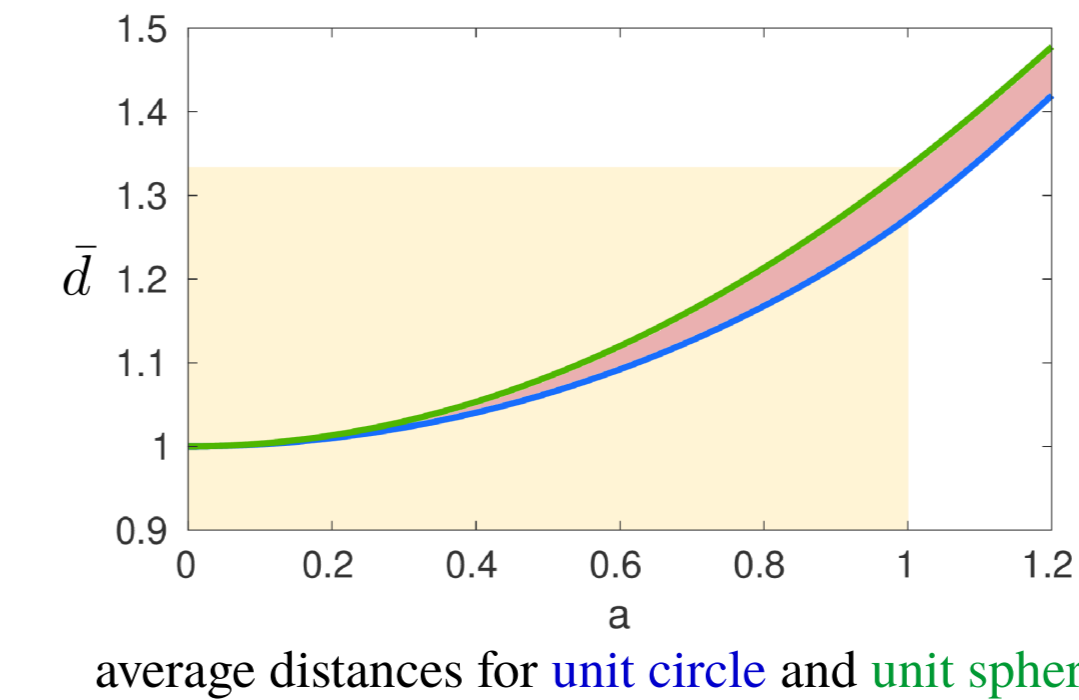
### Unit Sphere $S^2$

$$E_s = NM$$

$$\bar{d} = \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} 2 \sin \frac{\theta}{2} \sin \theta d\theta d\phi = \frac{4}{3}$$

$$E_a = \bar{d}NM = \frac{4NM}{3}$$

$$\bar{d}(a) = \frac{1}{4\pi} \int_0^\pi \int_0^{2\pi} \sqrt{a^2 + 1 - 2a \cos \theta} \sin \theta d\theta d\phi = \begin{cases} \frac{4}{3} & \text{if } a = 0, \\ \frac{1}{2} \int_0^\pi \sin \theta d\theta & \text{if } a > 0. \end{cases} = \begin{cases} 1, & \text{if } a = 0, \\ \frac{(a^2 + 1 - 2a \cos \theta)^{3/2}}{6a} \Big|_0^\pi & \text{if } a > 0. \end{cases}$$



average distances for unit circle and unit sphere

### Unit Disk

$$\bar{d} = \frac{1}{\pi} \int_0^1 2r \int_0^\pi r d\theta dr = \int_0^1 2r^2 dr = \frac{2}{3}$$

$$\bar{d} = \frac{1}{\pi} \int_0^1 2r \int_0^1 2q \int_0^\pi \sqrt{r^2 + q^2 - 2rq \cos \theta} d\theta dq dr = \frac{8}{\pi} \int_0^1 r \int_0^1 q |q-r| E\left(\frac{\theta}{2} \mid -\frac{4rq}{(q-r)^2}\right) dq dr \approx 0.9062$$

$$\bar{d} = \frac{125}{45\pi} \approx 0.9054 \quad [6]$$

$$\bar{d}(a) = \frac{8}{\pi} \int_0^1 r \int_0^a \frac{q}{a^2} |q-r| E\left(\frac{\theta}{2} \mid -\frac{4rq}{(q-r)^2}\right) dq dr$$

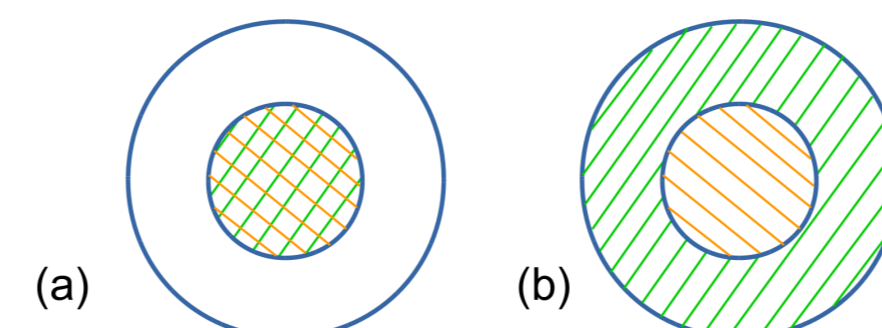
$$\bar{d} \approx \frac{8a}{\pi nm} \sum_{r=0}^1 r \sum_{q=0}^a \frac{q}{a^2} |q-r| E\left(\frac{\theta}{2} \mid -\frac{4rq}{(q-r)^2}\right) dq dr$$

Singularity

$$\bar{d} = \int_0^1 3r^2 \int_0^\pi r \frac{\sin \theta}{2} d\theta dr = \int_0^1 3r^3 dr = \frac{3}{4}$$

Alignment

$$\bar{d} = \int_0^1 3r^2 \int_0^1 3q^2 \int_0^\pi \frac{\sqrt{r^2 + q^2 - 2rq \cos \theta} \sin \theta}{2} d\theta dq dr$$

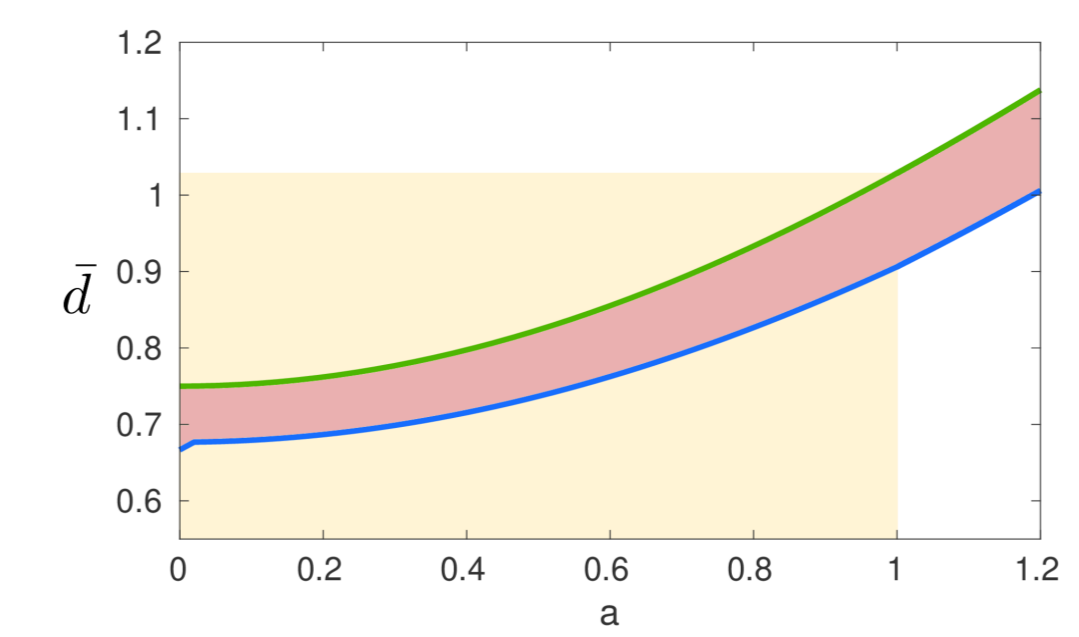


$$\int_0^\pi \frac{\sqrt{r^2 + q^2 - 2rq \cos \theta} \sin \theta}{2} d\theta \quad (r, q \geq 0)$$

$$\frac{(r+q)^3 - (r-q)^2|r-q|}{6rq} \quad (r \geq q) \quad \frac{6r^2q + 2q^3}{6rq}$$

$$\bar{d} = 2 \int_0^1 \int_0^1 \frac{9q^2 r^2 (6r^2q + 2q^3)}{6rq} dq dr = \frac{36}{35}$$

### Unit Ball



average distances for unit disk and unit ball

General Case

$$\bar{d}(a) = \int_0^1 3r^2 \int_0^a \frac{3q^2}{a^3} \int_0^\pi \frac{\sqrt{r^2 + q^2 - 2rq \cos \theta} \sin \theta}{2} d\theta dq dr$$

$$\bar{d}(a) = d_a(a) + d_{a,1}(a)$$

$$d_a(a) = 2 \int_0^a \int_0^a \frac{9q^2 r^2 (6r^2q + 2q^3)}{a^3 6rq} dq dr = \frac{36}{35} a^4$$

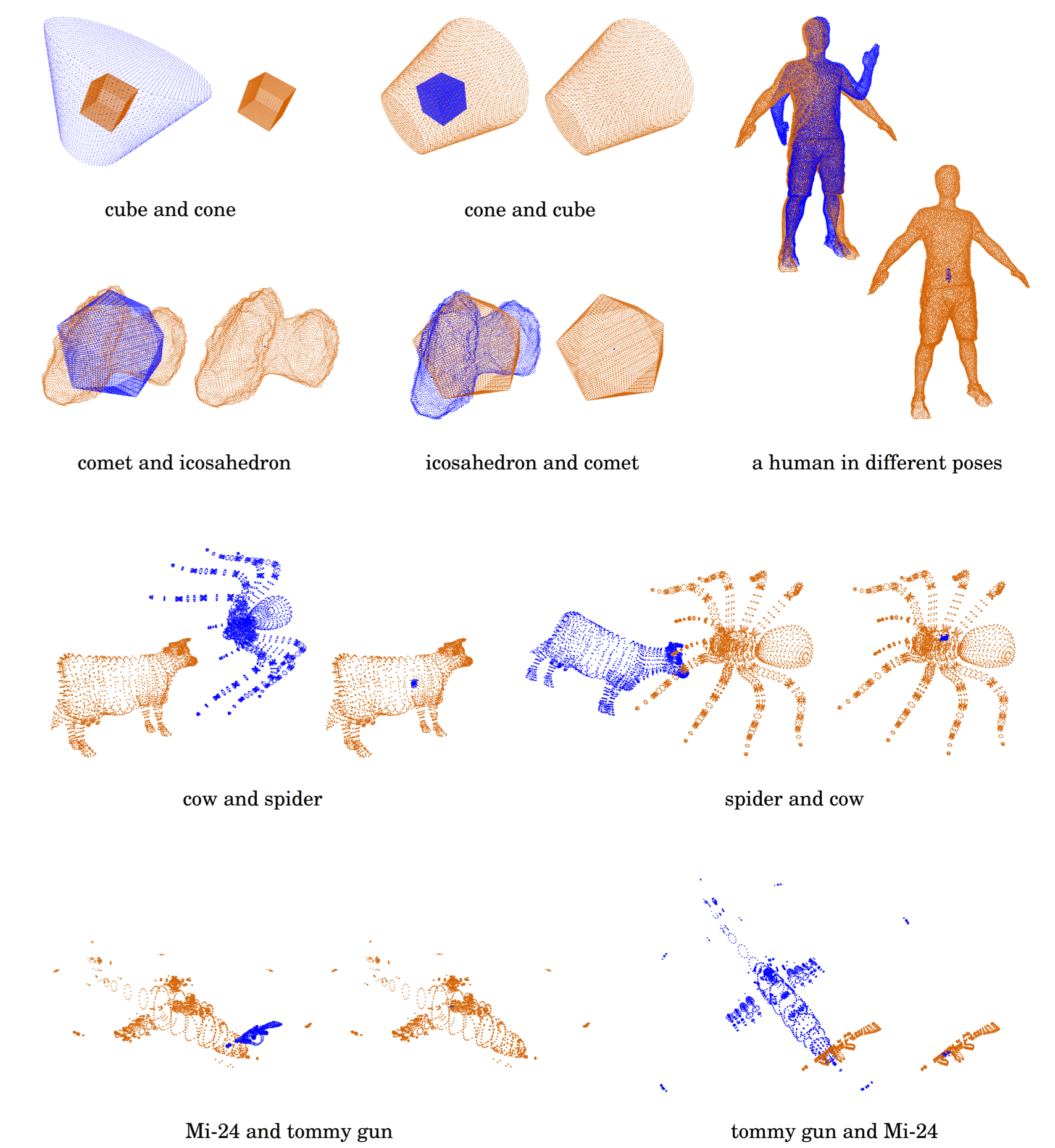
$$d_{a,1}(a) = \int_a^1 \int_0^a \frac{3rq}{2a^3} (6r^2q + 2q^3) dq dr = \frac{3}{5} r (5r^2 + a^2) dr = -\frac{21a^4 - 6a^2 - 15}{20}$$

$$\bar{d}(a) = \frac{36}{35} a^4 - \frac{21a^4 - 6a^2 - 15}{20} = \frac{26\frac{1}{4} + 10\frac{1}{2}a^2 - \frac{3}{4}a^4}{35}$$

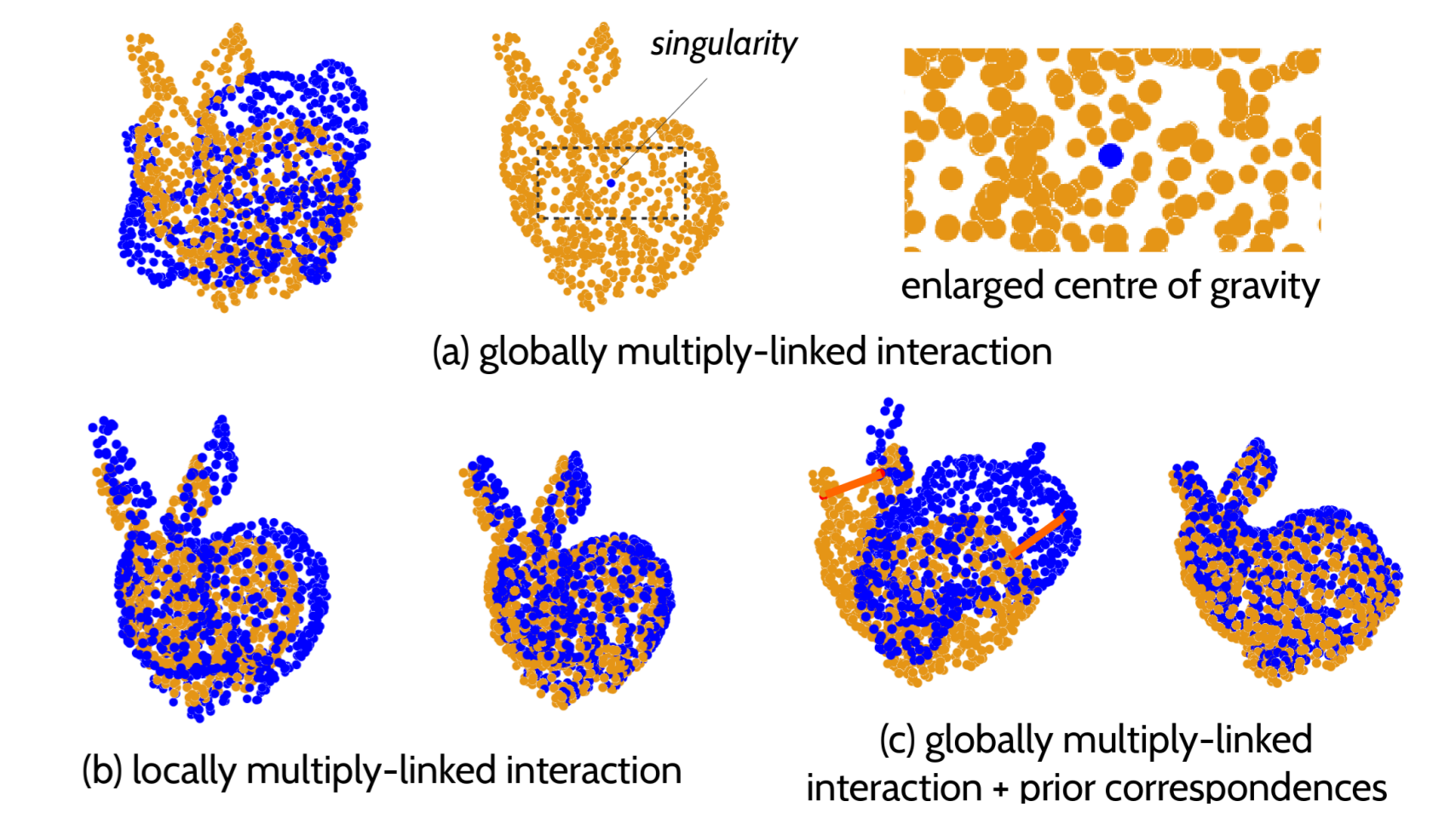
$$\bar{d}(0) = \frac{3}{4}$$

$$\bar{d}(1) = \frac{36}{35}$$

## OTHER SHAPES (TESTS WITH REAL DATA)



Selected registration results by a 7 DoF version of BH-RGA [8] with different shape combinations. For every pair of shapes, initialisation is on the left, and the alignment is on the right.



If template scaling is allowed, globally multiply-linked gravitational point set alignment results in a singularity (a). To remedy scaling resolution, point interactions can be restricted to local neighbourhoods (b) or prior matches with high masses can be embedded (c).

## WHAT IS A SCALE SINGULARITY?

- A singularity is a state when a template collapses to a single point, and scale  $\rightarrow 0$ .
- A singularity can arise in 7DoF point set alignment approaches due to numerical reasons (e.g., it can be observed in CPD [2]).
- In gravitational approaches, a singularity arises because of the underlying form of the gravitational potential energy (GPE) functional. We found that the GPE of a singularity is always smaller than the GPE of the optimal alignment under allowed scaling, and there are no equienergetic states with singularities.

## ELLIPTIC INTEGRALS

... of the second kind in the Legendre form are integrals of type

$$E(\phi, k) = \int_0^\phi \sqrt{1 - k \sin^2 \theta} d\theta$$

- Elliptic integrals arise in the study of the arc length problem for ellipses.
- As a rule, they cannot be simplified and analytically evaluated.

## APPENDIX

$$\star \lim_{a \rightarrow 1} \frac{1}{\pi} \int_0^\pi \sqrt{a^2 + 1 - 2a \cos \theta} d\theta = \frac{4}{\pi} \quad \blacklozenge \lim_{a \rightarrow 0} \frac{1}{2} \int_0^\pi \sqrt{1 + a^2 + 2a \cos \theta} \sin \theta d\theta = 1$$

$$\int \sqrt{a^2 + 1 - 2a \cos \theta} d\theta = \int \sqrt{a^2 + 1 + \frac{4a(1 - \cos \theta)}{2} - 2a \cos \theta} = \int \sqrt{4a \sin^2 \frac{\theta}{2} + \frac{(a-1)^4}{(a-1)^2}} d\theta = \sqrt{(a-1)^2} \int \sqrt{\frac{4a \sin^2 \frac{\theta}{2}}{(a-1)^2} + 1} d\theta = 2|a-1| E\left(\frac{\theta}{2} \mid -\frac{4a}{(a-1)^2}\right)$$

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